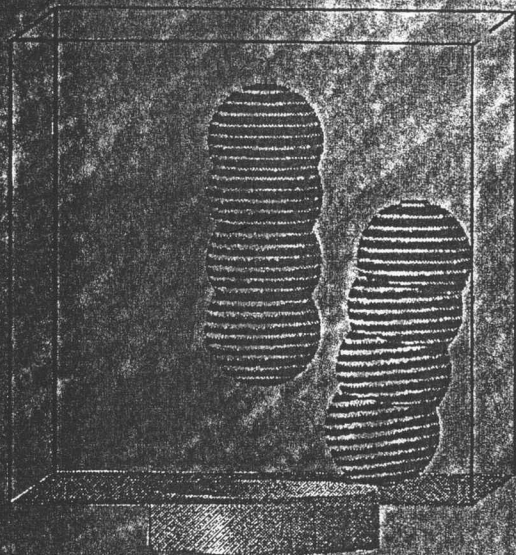
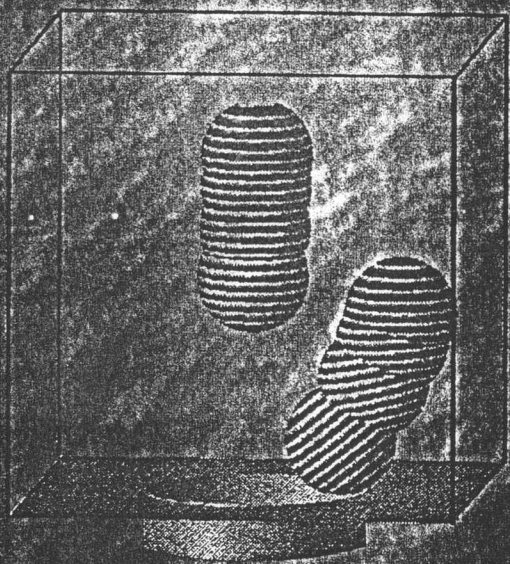


Boundary Element Technology XIII

C.S. Chen, C.A. Brebbia and D.W. Pepper
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Analysis of three-dimensional potential problems using the complex variable boundary element method

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Abstract

The Complex Variable Boundary Element Method, or CVBEM, is a two-dimensional (2D) potential problem numerical approximation technique. In this paper, the application of the CVBEM towards improving the numerical accuracy of three-dimensional (3D) numerical methods, in solving 3D potential problems, is introduced. The provided theoretical approach can be directly extended to other 2D numerical techniques.

I. Mathematical formulation

Let Ω be a three-dimensional (3D) domain, with boundary Γ , with coordinates given for an arbitrary point by (τ_1, τ_2, τ_3) .

Consider the 3D potential problem $\nabla^2\phi = f(\tau_1, \tau_2, \tau_3)$ on $\Omega \cup \Gamma$ with boundary conditions of the Dirichlet type (BCs) defined on Γ by the function $\phi_b(\tau_1, \tau_2, \tau_3)$. For a selected coordinate value $\tau_3 = \tau'_3$, the corresponding boundary is the boundary of a 2D "slice", $\Gamma(\tau'_3)$, denoted hereafter as Γ' , and the corresponding domain is the interior of the "slice", $\Omega(\tau'_3)$, denoted hereafter as Ω' , with boundary conditions on Γ' given by $\phi_b(\tau_1, \tau_2)$. Similarly, for $\tau_3 = \tau'_3$, $f(\tau_1, \tau_2, \tau'_3)$; $f(\tau_1, \tau_2)$. On the domain Ω' , which is subset of Ω , we have

can be chosen in an infinite number of directions. Our goal is to approximate $\frac{\partial^2 \phi}{\partial \tau_3^2}$ as closely as possible by $\frac{\partial^2 \phi^P}{\partial \tau_3^2}$.

III. Approximation of $\frac{\partial^2 \phi}{\partial \tau_3^2}$ by $\frac{\partial^2 \phi^P}{\partial \tau_3^2}$

A second order finite difference approximation of both terms in the equality setting $\frac{\partial^2 \phi}{\partial \tau_3^2} = \frac{\partial^2 \phi^P}{\partial \tau_3^2}$ is

$$\frac{\phi_{i+1} - 2\phi_i + \phi_{i-1}}{\Delta s^2} = \frac{\phi_{i+1}^P - 2\phi_i^P + \phi_{i-1}^P}{\Delta s^2} \quad (9)$$

where s is a tangential coordinate along the 3D boundary, Γ , and is chosen to include a nonzero τ_3 component (so as to be outside of the selected 2D slice), and i is an evaluation point on Γ , where n evaluation points are used for numerical integration purposes.

Then, at evaluation point i , we set $\phi_{i+1} - 2\phi_i + \phi_{i-1} = \eta_i = \phi_{i+1}^P - 2\phi_i^P + \phi_{i-1}^P$, or in terms of matrices and vectors, for n evaluation points on Γ to analyze $\frac{\partial^2 \phi}{\partial \tau_3^2}$,

$$\langle \eta \rangle_{n \times 1} = \begin{bmatrix} | \\ | \\ | \\ | \\ | \\ | \\ | \\ | \\ | \\ | \end{bmatrix} A \begin{bmatrix} | \\ | \\ | \\ | \\ | \\ | \\ | \\ | \\ | \\ | \end{bmatrix} = [A] \langle \phi^P \rangle \quad (10)$$

where A is a square $n \times n$ matrix; and $\langle \phi^P \rangle$ is a $n \times 1$ column vector.

But $\phi_i^P = \sum_{j=1}^N j g_j(p_i)$, such that the coordinate of evaluation point p_i is

$$(\tau_1, \tau_2, \tau_3)_i \quad (11)$$

$$\nabla^2\phi = \frac{\partial^2\phi}{\partial\tau_1^2} + \frac{\partial^2\phi}{\partial\tau_2^2} + \frac{\partial^2\phi}{\partial\tau_3^2} \quad (1)$$

which on $\Omega' \cup \Gamma$ reduces to,

$$\frac{\partial^2\phi}{\partial\tau_1^2} + \frac{\partial^2\phi}{\partial\tau_2^2} = f(\tau_1, \tau_2) - \frac{\partial^2\phi}{\partial\tau_3^2}, \text{ with BCs } \phi_b(\tau_1, \tau_2). \quad (2)$$

Let functions ϕ_f^P and ϕ^P satisfy

$$\frac{\partial^2\phi_f^P}{\partial\tau_1^2} + \frac{\partial^2\phi_f^P}{\partial\tau_2^2} = f(\tau_1, \tau_2) \quad (3)$$

and

$$\frac{\partial^2\phi^P}{\partial\tau_1^2} + \frac{\partial^2\phi^P}{\partial\tau_2^2} = -\frac{\partial^2\phi}{\partial\tau_3^2} \quad (4)$$

Then, on $\Omega' \cup \Gamma$, we can solve the governing 3D partial differential equation (PDE) by first solving the 2D PDE,

$$\frac{\partial^2\hat{\phi}}{\partial\tau_1^2} + \frac{\partial^2\hat{\phi}}{\partial\tau_2^2} = 0, \text{ with BCs } \hat{\phi}_b = \phi_b - \phi_f^P - \phi^P \text{ on } \Gamma'. \quad (5)$$

Then, the 3D approximation in $\Omega' \cup \Gamma$, is $\hat{\Phi} = \hat{\phi} + \phi_f^P + \phi^P$. (6)

II. Numerical modeling approach

This paper's approach to solving (2) is to use a 3D function ϕ^P of the form

$$\phi^P = \sum_{j=1}^{N3} c_j g_j \quad (7)$$

where the g_j are mutually independent functions such that the 3D Laplacian $\nabla^2 g_j = 0$. The notation, $N3$, is the number of 3D nodes that enclose $\Omega \cup \Gamma$. We select c_j so that ϵ_1 is minimized in the usual residual error sense, where

$$\epsilon_1 = \left| \phi_{ss}^P - \phi_{ss} \right|_{\Gamma'} \quad (8)$$

where ϕ_{ss} is the tangential second partial derivative of ϕ along Ω . Note that s

But from (11), the column vector of values ϕ_i^p is given by,

$$\langle \phi^p \rangle_{n \times 1} = \begin{bmatrix} g_1(p_1) & g_2(p_1) & \cdots & g_N(p_1) \\ g_1(p_2) & g_2(p_2) & \cdots & g_N(p_2) \\ \cdot & \cdot & & \cdot \\ \cdot & \cdot & & \cdot \\ \cdot & \cdot & & \cdot \\ g_1(p_n) & g_2(p_n) & \cdots & g_N(p_n) \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \\ \cdot \\ \cdot \\ \cdot \\ c_N \end{bmatrix} \quad (12)$$

$$= [g(p)] \langle c \rangle$$

where $n \geq N$.

Then, combining (11) and (12),

$$\langle \eta \rangle_{n \times 1} = [A]_{n \times n} [g(p)]_{n \times N} \langle c \rangle_{N \times 1} \quad (13)$$

where matrices $[A]$ and $[g(p)]$ are completely determined by the above definitions, and $\langle c \rangle$ is to be determined by a least-squares type solution to minimize the usual residual error norm. Note that $\langle \eta \rangle = [A] \langle \phi \rangle$ where $\langle \phi \rangle = (\phi_1, \phi_2, \dots, \phi_n)$. Then, equating matrix relationships,

$$[A] \langle \phi \rangle = [A] [g(p)] \langle c \rangle \quad (14)$$

or, for $[A]$ nonsingular,

$$\langle \phi \rangle = [g(p)] \langle c \rangle \quad (15)$$

which is the typical direct formulation for developing a 3D approximator. Thus, the standard 3D numerical technique also is a model for estimating $\frac{f^2 \phi}{f \tau_3^2}$.

Thus, a typical 3D solution of the Laplace equation can be directly used as ϕ^p .

IV. The complex variable boundary element method (CVBEM)

The CVBEM will be used to develop the 2D approximation, on $\Omega' \cup \Gamma'$ according to (5). Background on the CVBEM is provided in numerous texts and papers (see references).

V. The 3D model particular solution, ϕ^P

Given ϕ^P , and ϕ_f^P , we now develop new BCs along Γ' to be used in solving for the 2D approximator $\hat{\phi}$ by the CVBEM. Note that these new BC values, on Γ' , is the discrepancy between $(\phi^P + \phi_f^P)$ and ϕ along Γ' .

VI. Vector space considerations

The ϕ^P approximator utilizes 3D basis functions of the form $g_j(\tau_1, \tau_2, \tau_3)$. The CVBEM utilizes 2D basis functions of the form $h_j(\tau_1, \tau_2)$. The linear operator is $L = \nabla^2$. Then, for any j , and independently,

$$Lg_j = 0 ; Lh_j = 0. (16)$$

Recall, $L\phi_f^P = f$. Then the basis developed in this numerical approach is $B = \{g_j; h_j; \phi_f^P\}$.

VII. Discussion

In general, accuracy in 3D problems is increased by adding 3D basis functions, which increases modeling complexity. In our new case, we get another chance towards reducing modeling error by, in effect, adding additional basis functions in a 2D subspace of the 3D domain. These additional 2D basis functions can only help in reducing the overall 3D modeling error, in a least squares sense, although such improvements are seen on a 2D "slice" by "slice" basis.

VII. Conclusions

The Complex Variable Boundary Element Method, or CVBEM, is a two-dimensional (2D) potential problem numerical approximation technique. In this paper, the application of the CVBEM towards improving the numerical accuracy of three-dimensional (3D) numerical methods, in solving 3D potential problems, is introduced. The theoretical approach provided can be directly extended to other 2D numerical techniques.

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