

# EXPANDING THE CVBEM ANALOG INTO A SERIES

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In the past, the CVBEM (complex variable boundary element method) has been approached as a collocation problem or by least squares. In this work, the CVBEM analog is redeveloped, for the first time, as a series expansion of nodal point functions with unknown nodal points values as the coefficients. This series expansion provides further insight into the theoretical and approximation aspects of the CVBEM.

*Key words:* CVBEM, series expansion, boundary elements, complex variables, complex variable boundary element method.

## Introduction

The complex variable boundary element method, or CVBEM, has gained increased use in approximating two-dimensional potential problems since its inception nearly ten years ago.<sup>1</sup> The CVBEM has been extended to include collocation techniques, least-squares minimization, and use of singular approximation functions. In this note, the CVBEM numerical analog is reformulated into a new expansion of nodal point approximation functions such that the several techniques described previously can be unified. With this new formulation, the CVBEM is now described as a series expansion, developed from the Cauchy integral.

## Formulation

The CVBEM is developed from numerically approximating the Cauchy contour integral<sup>1</sup>

$$\omega(z) = \frac{1}{2\pi i} \int_{\Gamma} \frac{\omega(\zeta) d\zeta}{\zeta - z} \quad (1)$$

where  $\Gamma$  is a simple closed contour enclosing a simply connected domain  $\Omega$ ;  $\omega(z)$  is an analytic function on  $\Omega \cup \Gamma$ ;  $\zeta$  is an integration variable;  $i = \sqrt{-1}$ . It is assumed in (1) that  $\Gamma$  is a polygon with  $v$  vertices, and values of  $\omega(\zeta)$  are given on  $\Gamma$  such as to provide a well-defined mixed boundary value potential problem.

Contour  $\Gamma$  is discretized into  $m$  boundary elements,  $\Gamma_j$ , by  $m$  nodal points such that  $m > v$ , and a node is located at each vertex of  $\Gamma$ .<sup>1</sup> A piecewise linear spline function  $N_j(\zeta)$  is defined for each node  $j$  (located coordinate  $z_j \in \Gamma$ ) by

$$N_j(\zeta) = \begin{cases} 0, & \zeta \notin \Gamma_j \cup \Gamma_{j-1} \\ (\zeta - z_j)/(z_j - z_{j-1}), & \zeta \in \Gamma_{j-1} \\ (z_{j+1} - \zeta)/(z_{j+1} - z_j), & \zeta \in \Gamma_j \end{cases} \quad (2)$$

where  $z_{m+1} \equiv z_1$  appropriately.

A global trial function, based upon the sum of nodal contributions is

$$G(\zeta) = \sum_{j=1}^m N_j(\zeta) \omega_j \quad (3)$$

where  $G(\zeta)$  is a continuous piecewise linear function on  $\Gamma$ ;  $\omega_j$  are nodal point values of a potential function  $\phi(z)$  and its conjugate  $\psi(z)$ , where  $\omega_j = \phi_j + i\psi_j$ . (Generally, nodal points are located on  $\Gamma$  such that the given boundary conditions are matched on  $\Gamma$  by the global trial function in (3). Higher order spline functions may be used in (3), or an increase in nodal point density, in order to obtain a match of the known values of  $\phi(z)$  on  $\Gamma$ ).

The CVBEM approximator, for linear trial functions (higher order polynomial splines may be used directly in this development) is

$$\hat{\omega}(z) = \frac{1}{2\pi i} \int_{\Gamma} \frac{G(\zeta) d\zeta}{\zeta - z}, \quad z \in \Omega \quad (4)$$

where  $\hat{\omega}(z)$  is analytic in  $\Omega$  and continuous on  $\Gamma$ . Note that  $\hat{\omega}(z) = \hat{\phi}(z) + i\hat{\psi}(z)$ , and  $\nabla^2 \hat{\phi}(z) = 0$ ,  $\nabla^2 \hat{\psi}(z) = 0$ .

## Series expansion of CVBEM analog

The formulation of (4) is now expanded by substituting (3) into (4) by

$$\hat{\omega}(z) = \frac{1}{2\pi i} \int_{\Gamma} \sum_{j=1}^m \frac{\omega_j N_j(\zeta) d\zeta}{\zeta - z} \quad (5)$$

$$= \sum_{j=1}^m \frac{1}{2\pi i} \int_{\Gamma_{j-1} \cup \Gamma_j} \frac{\omega_j N_j(\zeta) d\zeta}{\zeta - z} \quad (6)$$

where in (6),  $\Gamma_0 \equiv \Gamma_m$  appropriately.

Rewriting (6),

$$\hat{\omega}(z) = \frac{1}{2\pi i} \sum_{j=1}^m \omega_j \int_{\Gamma_{j-1} \cup \Gamma_j} \frac{N_j(\zeta) d\zeta}{\zeta - z} \quad (7)$$

Integrating (7), the series expansion is developed by

$$\begin{aligned} \hat{\omega}(z) &= \sum_{j=1}^m \omega_j \frac{1}{2\pi i} \left[ \left( \frac{z - z_{j-1}}{z_j - z_{j-1}} \right) \right. \\ &\quad \left. (\ln(z_j - z) - \ln(z_j - z)) + \left( \frac{z_{j+1} - z}{z_{j+1} - z_j} \right) \right. \\ &\quad \left. (\ln(z_j - z) - \ln(z_{j+1} - z)) \right] \\ &= \sum_{j=1}^m \omega_j H_j(z) \end{aligned} \quad (8)$$

where in (9),  $H_j(z)$  follows immediately from (8).

Note that in (8) and (9), the limit of  $\hat{\omega}(z)$ , as point  $z \in \Omega$  approaches any nodal coordinate  $z_j \in \Omega$  is well defined; this limiting value leads to the collocation techniques used in CVBEM.<sup>1</sup> Also, the complex logarithm branch-cut issues are avoided by the expansion of (8).

The least squares techniques used in the CVBEM<sup>3</sup> can be seen in using (9) by first considering the known and the unknown nodal values given respectively by  $\xi_p^k$ ,  $\xi_p^u$ . Letting  $\psi$  be a characteristic type function assuming the values 1 or  $i$ , for  $\xi$  being real or imaginary, respectively, then (9) is written as the series expansion of nodal point functions by

$$\hat{\omega}(z) = \sum_{p=1}^{nu} \psi_p^u \xi_p^u H_p(z) + \sum_{q=1}^{nk} \psi_q^k \xi_q^k H_q(z) \quad (10)$$

where  $\psi_p^u$  and  $\psi_q^k$  are the complex-valued characteristic type functions associated with the unknown and known nodal values, respectively;  $nu$  and  $nk$  are the number of unknown and known nodal values, respectively; and  $nu + nk = 2m$ .

In (10), the  $\xi_q^k$  are known by the boundary conditions of the problem, and the unknown nodal values,  $\xi_p^u$ , are to be estimated by the  $\xi_p^u$ .

From (10)

$$\hat{\omega}(z) = f^u(z) = f^k(z) \quad (11)$$

where  $f^k(z)$  is a known function of  $z$ , and

$$f^u(z) = \sum_{p=1}^{nu} \psi_p^u \xi_p^u H_p(z) \quad (12)$$

is a function of the  $nu$  unknown nodal value estimates  $\{\xi_p^u, p=1, 2, \dots, nu\}$  and the complex variable  $z \in \Omega$ . Equations (10) through (12) present a new series expansion for the CVBEM.

The CVBEM formulation can now be written as a minimization of a norm

$$\| f^u(z) - (\omega(z) - f^k(z)) \| \quad (13)$$

on the problem boundary,  $\Gamma$ , with respect to the given boundary condition values. With a match of the boundary conditions, given on  $\Gamma$ , by the selected spline functions and  $\xi_q^k$  values, there is no approximation error associated with the complex function  $f^k(z)$ .

Thus, the CVBEM approximation analog is to minimize the error in fitting, in a least squares sense (or other norm), the difference in boundary values of  $f^u(z)$  and  $(\omega(z) - f^k(z))$  on  $\Gamma$ , where boundary condition values are given.

## Conclusions

The CVBEM analog is, for the first time, redeveloped into a series expansion of nodal point functions. From the series expansion, the well-known collocation and least-squares formulations currently in use, can be seen to be applications of different norms in minimizing a fitting of the CVBEM approximation function,  $f^u(z)$ , to the boundary conditions of the problem.

## References

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